

ELECTRON-PHOTON INTERACTION

Photon Absorption Rate

$$V(\vec{R}, t) = \frac{eE_0}{im_0\omega} \left[e^{i(\vec{Q}\cdot\vec{R}-\omega t)} - e^{-i(\vec{Q}\cdot\vec{R}-\omega t)} \right] (\vec{e} \cdot \hat{p})$$

- The transition rate from a state i to another j due to the absorption of a photon

$$W_{ji} = \frac{2\pi}{\hbar} \left(\frac{eE_0}{m_0\omega} \right)^2 \left| \langle j | \vec{e} \cdot \hat{p} | i \rangle \right|^2 \delta(E_j - E_i - \hbar\omega)$$

$$\vec{e} \cdot \hat{p} = -i\hbar \left(e_x \frac{\partial}{\partial x} + e_y \frac{\partial}{\partial y} + e_z \frac{\partial}{\partial z} \right)$$

Power Absorbed

- Power absorbed in $i \rightarrow j$ transition = $W_{ji} \hbar \omega$

- Total Power absorbed, $P_+ = \hbar \omega \sum_{i,j} W_{ji}$

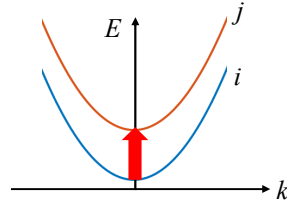
- Include the Fermi distribution:

$$P_+ = \hbar \omega \sum_{i,j} W_{ji} f(E_i) [1 - f(E_j)]$$

$f(E_i)$: Initial state is filled

$1 - f(E_j)$: Final state is empty

- Consider the spin degeneracy: $P_+ = 2\hbar \omega \sum_{i,j} W_{ji} f(E_i) [1 - f(E_j)]$

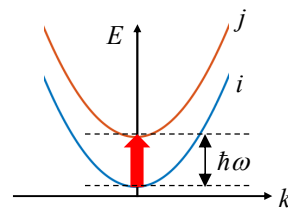


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Power Absorption and Emission

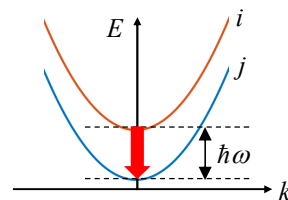
- Absorption:

$$P_+ = \frac{2\pi}{\hbar} \hbar \omega \left(\frac{eE_0}{m_0\omega} \right)^2 2 \sum_{i,j} |\langle j | \vec{e} \cdot \hat{p} | i \rangle|^2 \times f(E_i) [1 - f(E_j)] \delta(E_j - E_i - \hbar\omega)$$



- Emission:

$$P_- = -\frac{2\pi}{\hbar} \hbar \omega \left(\frac{eE_0}{m_0\omega} \right)^2 2 \sum_{i,j} |\langle j | \vec{e} \cdot \hat{p} | i \rangle|^2 \times f(E_i) [1 - f(E_j)] \delta(E_j - E_i + \hbar\omega)$$



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Net Power Absorption

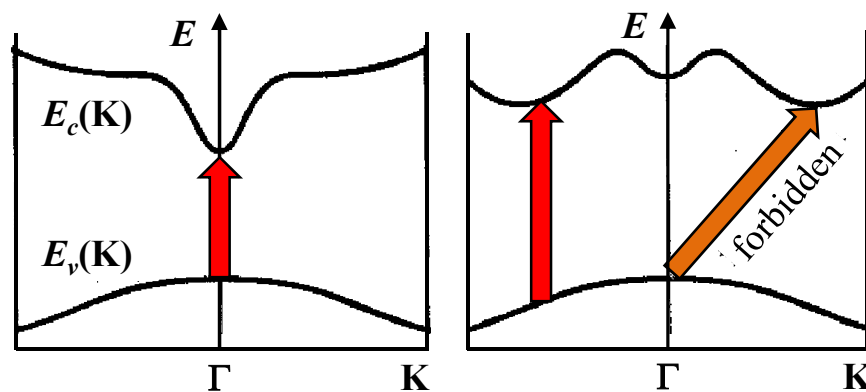
$$f(E_i)[1-f(E_j)] - f(E_j)[1-f(E_i)] = f(E_i) - f(E_j)$$

$$P = \frac{2\pi}{\hbar} \hbar\omega \left(\frac{eE_0}{m_0\omega} \right)^2 \sum_{i,j} |\langle j | \vec{e} \cdot \hat{p} | i \rangle|^2 \times [f(E_i) - f(E_j)] \delta(E_j - E_i - \hbar\omega)$$

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Optical Absorption

- Band gaps can be measured by optical absorption.

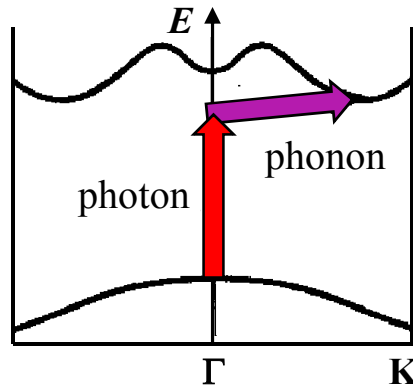


$$E_f = E_i + \hbar\omega, \quad \vec{K}_f = \vec{K}_i + \vec{Q}$$

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Optical Absorption

- Combination of photon and phonon transitions across the bandgap.
- Direct bandgap materials, e.g., GaAs, are preferable to indirect bandgap materials, e.g., Si, for making lasers or detectors.



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Interband Absorption

- **Classical Approach:**

- Consider a light wave -

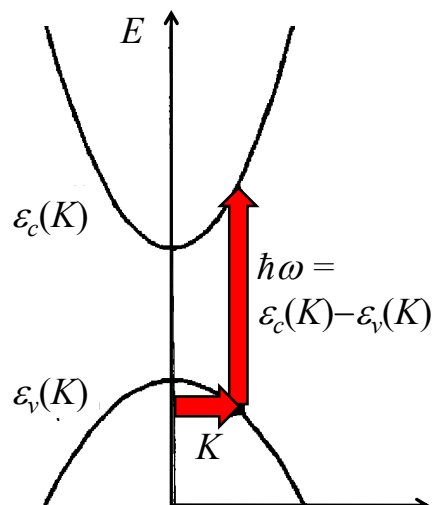
$$\vec{E}(\vec{R}, t) = 2\bar{e}E_0 \cos(\vec{Q} \cdot \vec{R} - \omega t)$$

- \vec{E} induces a current with in-phase component -

$$\vec{J} = \sigma_1 \vec{E}$$

- Energy dissipation rate -

$$\vec{J} \cdot \vec{E}$$



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Classical Approach

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- Consider a light wave -

$$\vec{E}(\vec{R}, t) = 2\bar{e}E_0 \cos(\vec{Q} \cdot \vec{R} - \omega t)$$

- \vec{E} induces a current with in-phase component -

$$\vec{J} = \sigma_1 \vec{E}$$

- Energy dissipation rate -

$$\vec{J} \cdot \vec{E}$$

- **Total power dissipated:**

$$\begin{aligned} P &= \Omega \frac{1}{t} \int_0^t \vec{J} \cdot \vec{E} dt' \\ &= \frac{4\sigma_1 \Omega E_0^2}{t} \int_0^t \cos^2(\vec{Q} \cdot \vec{R} - \omega t') dt' \\ &= \frac{2\sigma_1 \Omega E_0^2}{t} \int_0^t [1 + \cos 2(\vec{Q} \cdot \vec{R} - \omega t')] dt' \\ &= 2\sigma_1 \Omega E_0^2 \end{aligned}$$

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Conductivity

- Classical approach:

$$P = 2\sigma_1 \Omega E_0^2$$

- Quantum mechanical approach:

$$\begin{aligned} P &= \frac{2\pi}{\hbar} \hbar \omega \left(\frac{eE_0}{m_0 \omega} \right)^2 2 \sum_{i,j} |\langle j | \vec{e} \cdot \hat{p} | i \rangle|^2 \\ &\quad \times [f(E_i) - f(E_j)] \delta(E_j - E_i - \hbar \omega) \end{aligned}$$

$$\sigma_1(\omega) = \frac{\pi e^2}{m_0^2 \omega} \frac{2}{\Omega} \sum_{i,j} |\langle j | \vec{e} \cdot \hat{p} | i \rangle|^2 [f(E_i) - f(E_j)] \delta(E_j - E_i - \hbar \omega)$$

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