

ELECTROMAGNETIC WAVE PROPAGATION

1

Lossless Dielectric

$$\sigma \ll \omega \epsilon$$

$$\sigma \ll 0, \epsilon = \epsilon_0 \epsilon_r, \mu = \mu_0 \mu_r$$

$$\alpha = \omega \sqrt{\frac{\mu \epsilon}{2} \left[\sqrt{1 + \left[\frac{\sigma}{\omega \epsilon} \right]^2} - 1 \right]} = 0 \quad \beta = \omega \sqrt{\frac{\mu \epsilon}{2} \left[\sqrt{1 + \left[\frac{\sigma}{\omega \epsilon} \right]^2} + 1 \right]} = \omega \sqrt{\mu \epsilon}$$

$$\eta = j \frac{\omega \mu}{\gamma} = \sqrt{\frac{\mu}{\epsilon}} \angle 0^\circ$$

E and H are in time phase with each other.

$$\vec{E}(z, t) = E_0 e^{-\alpha z} \cos(\omega t - \beta z) \hat{a}_x = E_0 \cos(\omega t - \beta z) \hat{a}_x$$

$$\vec{H}(z, t) = (E_0 / \eta) e^{-\alpha z} \cos(\omega t - \beta z) \hat{a}_x = \sqrt{(\epsilon / \mu)} E_0 \cos(\omega t - \beta z) \hat{a}_x$$

2

Free Space: Intrinsic Impedance

$$\sigma = 0, \epsilon = \epsilon_0, \mu = \mu_0$$

$$\alpha = 0, \beta = \omega\sqrt{\mu_0\epsilon_0} = \frac{\omega}{c}$$

$$\eta_0 = \sqrt{\frac{\mu_0}{\epsilon_0}} = \sqrt{\frac{4\pi \times 10^{-7}}{8.854 \times 10^{-12}}} = 120\pi \approx 377 \Omega$$

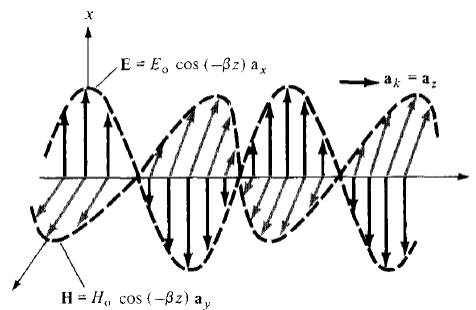
3

EM Waves in Free Space

$$\vec{E} = E_0 \cos(\omega t - \beta z) \hat{a}_x$$

$$\vec{H} = H_0 \cos(\omega t - \beta z) \hat{a}_y = \frac{E_0}{\eta_0} \cos(\omega t - \beta z) \hat{a}_y$$

$$\hat{a}_E \times \hat{a}_H = \hat{a}_k \rightarrow \text{TEM waves}$$



4

EM Waves in Good Conductors

$$\sigma \gg \omega, \epsilon = \epsilon_0 \epsilon_r, \mu = \mu_0 \mu_r$$

$$\alpha = \omega \sqrt{\frac{\mu\epsilon}{2} \left[\sqrt{1 + \left[\frac{\sigma}{\omega\epsilon} \right]^2} - 1 \right]} \quad \beta = \omega \sqrt{\frac{\mu\epsilon}{2} \left[\sqrt{1 + \left[\frac{\sigma}{\omega\epsilon} \right]^2} + 1 \right]}$$

$$\alpha = \beta = \sqrt{\frac{\omega\mu\sigma}{2}} = \sqrt{\pi f \mu \sigma}$$

$$\eta = \sqrt{\frac{j\omega\mu}{\sigma + j\omega\epsilon}} = \sqrt{\frac{\omega\mu}{\sigma}} \angle 45^\circ$$

E leads H by 45°.

5

EM Waves in Good Conductors

$$\vec{E} = E_0 e^{-\alpha z} \cos(\omega t - \beta z) \hat{a}_x$$

$$\vec{H} = \frac{E_0}{\sqrt{\omega\mu/\sigma}} e^{-\alpha z} \cos(\omega t - \beta z - 45^\circ) \hat{a}_y$$

Amplitude is attenuated by the factor $e^{-\alpha z}$.

6

ELECTROMAGNETIC WAVES AND CURRENT IN CONDUCTORS

7

EM Waves in Good Conductors

$$\sigma \gg \omega \epsilon$$

$$\alpha = \omega \sqrt{\frac{\mu \epsilon}{2} \left[\sqrt{1 + \left[\frac{\sigma}{\omega \epsilon} \right]^2} - 1 \right]} \quad \beta = \omega \sqrt{\frac{\mu \epsilon}{2} \left[\sqrt{1 + \left[\frac{\sigma}{\omega \epsilon} \right]^2} + 1 \right]}$$

$$\alpha = \beta = \sqrt{\frac{\omega \mu \sigma}{2}} = \sqrt{\pi f \mu \sigma}$$

$$\eta = \sqrt{\frac{j \omega \mu}{\sigma + j \omega \epsilon}} = \sqrt{\frac{\omega \mu}{\sigma}} \angle 45^\circ$$

Electric field leads magnetic field by 45°.

8

EM Waves in Good Conductors

$$\vec{E} = E_0 e^{-\alpha z} \cos(\omega t - \beta z) \hat{a}_x$$

$$\vec{H} = \frac{E_0}{\sqrt{\omega\mu/\sigma}} e^{-\alpha z} \cos(\omega t - \beta z - 45^\circ) \hat{a}_y$$

Amplitude is attenuated by the factor $e^{-\alpha z}$.

9

Skin Depth

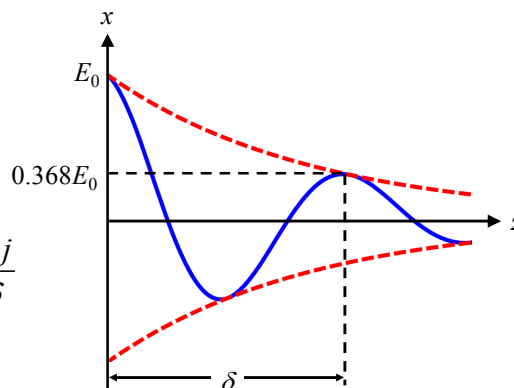
Skin depth, δ : Wave amplitude decreases by a factor e^{-1} ($\sim 37\%$).

$$E_0 e^{-\alpha\delta} = E_0 e^{-1}$$

$$\delta = \frac{1}{\alpha} \Rightarrow \delta = \sqrt{\frac{2}{\omega\mu\sigma}}$$

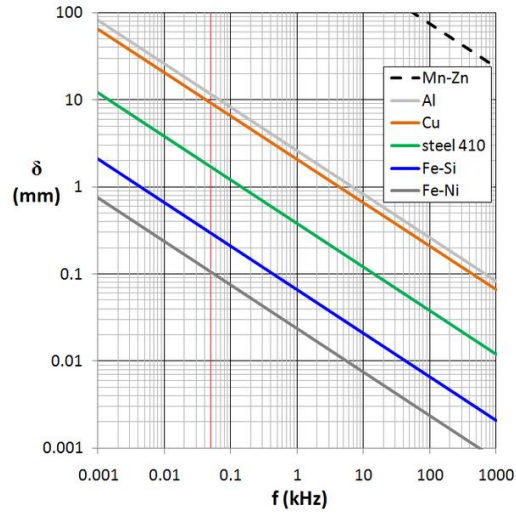
$$\eta = \sqrt{\frac{\omega\mu}{\sigma}} e^{j\pi/4} = \frac{\sqrt{2}}{\sigma\delta} e^{j\pi/4} = \frac{1+j}{\sigma\delta}$$

$$\vec{E} = E_0 e^{-z/\delta} \cos\left(\omega t - \frac{z}{\delta}\right) \hat{a}_x$$



10

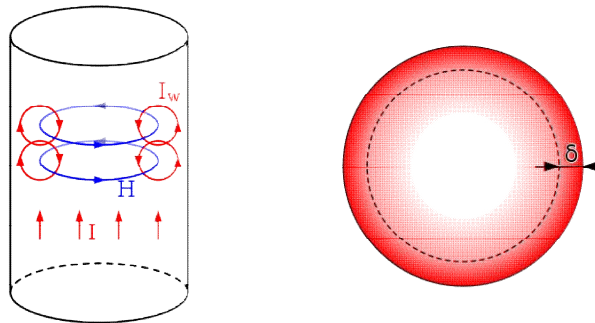
Skin Depth vs. Frequency



Skin effect is used as electromagnetic shielding!

11

Skin Effect

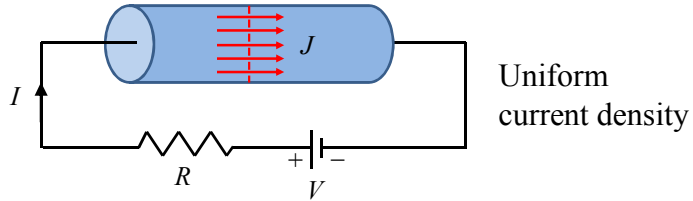


- Skin depth is due to the circulating eddy currents (arising from a changing H field) cancelling the current flow in the center of a conductor and reinforcing it in the skin.
- Distribution of current flow in a cylindrical conductor, shown in cross section. For alternating current, most (63%) of the electrical current flows between the surface and the skin depth, δ .

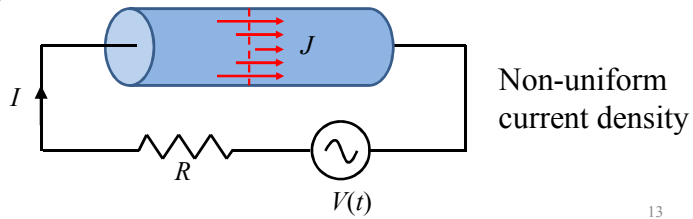
12

Current Flow in a Good Conductor

DC:



AC:



13

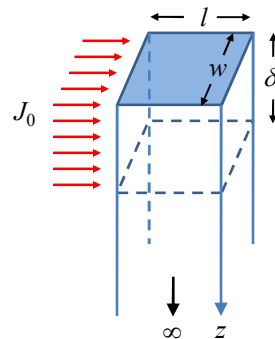
AC Resistance

$$R_{dc} = \frac{l}{\sigma S}$$

$$\eta = \frac{1}{\sigma\delta} + \frac{j}{\sigma\delta} = R_s + jX_s$$

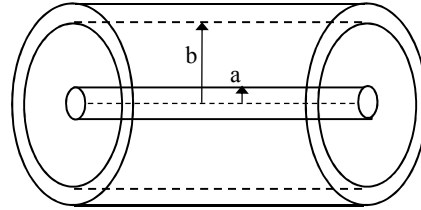
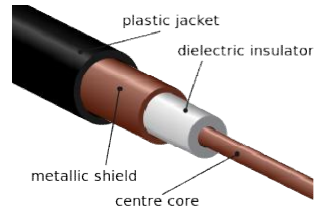
$$R_{ac} = \frac{l}{\sigma\delta w} = R_s \frac{l}{w}$$

R_{ac} is equivalent to R_{dc} of a plane conductor of length l and cross section $A = \delta w$.



14

Co-Axial Cable



The current flowing through the co-axial cable is concentrated within a thin layer on the outer surface of the inner conductor and on the inner surface of the outer conductor.



Equivalent inner conductor



Equivalent outer conductor