

# ELECTROMAGNETIC WAVES

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## *Electromagnetic Potentials*

**Electrostatics:**  $\nabla \times \vec{E} = 0 \Rightarrow \vec{E} = -\nabla V$

**What happens when the fields are time-varying?**

$$\nabla \times \vec{E} = -\frac{\partial B}{\partial t} = -\frac{\partial}{\partial t}(\nabla \times \vec{A}) \Rightarrow \nabla \times \left( \vec{E} + \frac{\partial \vec{A}}{\partial t} \right) = 0$$

Since the curl of the gradient of a scalar field is zero

$$\vec{E} + \frac{\partial \vec{A}}{\partial t} = -\nabla V$$

$$\vec{E} = -\nabla V - \frac{\partial \vec{A}}{\partial t}$$

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## *Time-Varying Potentials*

Maxwell's equation  $\rightarrow \nabla \cdot \vec{D} = \rho_v \Rightarrow \nabla \cdot \vec{E} = \frac{\rho_v}{\epsilon}$

Electromagnetic potential  $\rightarrow \vec{E} = -\nabla V - \frac{\partial \vec{A}}{\partial t}$

Taking divergence  $\rightarrow \nabla \cdot \vec{E} = -\nabla^2 V - \frac{\partial}{\partial t} (\nabla \cdot \vec{A}) = \frac{\rho_v}{\epsilon}$

Maxwell's equation  $\rightarrow \nabla \times \vec{H} = \vec{J} + \frac{\partial \vec{D}}{\partial t} \Rightarrow \nabla \times \vec{B} = \mu \vec{J} + \epsilon \mu \frac{\partial \vec{E}}{\partial t}$

$$\begin{aligned} \nabla \times \vec{B} &= \nabla \times \nabla \times \vec{A} = \mu \vec{J} + \epsilon \mu \frac{\partial}{\partial t} \left( -\nabla V - \frac{\partial \vec{A}}{\partial t} \right) \\ &= \mu \vec{J} - \mu \epsilon \nabla \left( \frac{\partial V}{\partial t} \right) - \mu \epsilon \frac{\partial^2 \vec{A}}{\partial t^2} \end{aligned}$$

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## *Time-Varying Potentials*

$$\nabla \times \nabla \times \vec{A} = \mu \vec{J} - \mu \epsilon \nabla \left( \frac{\partial V}{\partial t} \right) - \mu \epsilon \frac{\partial^2 \vec{A}}{\partial t^2}$$

Vector identity  $\rightarrow \nabla \times \nabla \times \vec{A} = \nabla(\nabla \cdot \vec{A}) - \nabla^2 \vec{A}$

$$\nabla^2 \vec{A} - \nabla(\nabla \cdot \vec{A}) = -\mu \vec{J} + \mu \epsilon \nabla \left( \frac{\partial V}{\partial t} \right) + \mu \epsilon \frac{\partial^2 \vec{A}}{\partial t^2}$$

Lorentz condition for potentials  $\rightarrow$

$$\nabla \cdot \vec{A} = -\mu \epsilon \frac{\partial V}{\partial t}$$

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## *Wave Equations*

$$\nabla^2 V - \mu\epsilon \frac{\partial^2 V}{\partial t^2} = -\frac{\rho_v}{\epsilon}$$

$$\nabla^2 \vec{A} - \mu\epsilon \frac{\partial^2 \vec{A}}{\partial t^2} = -\mu \vec{J}$$

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## *Maxwell's Equations*

*Free space with charges*

Differential form	Integral form
E-Gauss: $\nabla \cdot \epsilon_o \vec{E} = \rho$	$\oiint_S \epsilon_o \vec{E} \cdot d\vec{S} = \iiint_V \rho dV$
Faraday: $\nabla \times \vec{E} = -\frac{\partial}{\partial t} \mu_o \vec{H}$	$\oint_C \vec{E} \cdot d\vec{C} = \iint_S -\frac{\partial}{\partial t} \mu_o \vec{H} \cdot d\vec{S}$
H-Gauss: $\nabla \cdot \mu_o \vec{H} = 0$	$\oiint_S \mu_o \vec{H} \cdot d\vec{S} = 0$
Ampere: $\nabla \times \vec{H} = \vec{J} + \frac{\partial}{\partial t} \epsilon_o \vec{E}$	$\oint_C \vec{H} \cdot d\vec{C} = \iint_S (\vec{J} + \frac{\partial}{\partial t} \epsilon_o \vec{E}) \cdot d\vec{S}$

In statics, both time derivatives are unimportant, Maxwell's Equations split into decoupled electrostatic and magnetostatic equations.

In Electro-quasistatic (EQS) and magneto-quasistatic systems (MQS), one (but not both) time derivative becomes important.

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## Coupling of Electric and Magnetic Fields

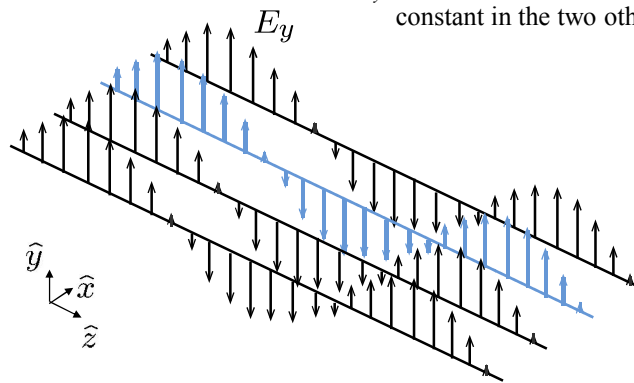
Maxwell's Equations couple the E and H fields:

$$\oint_C \vec{E} \cdot d\vec{l} = -\frac{d}{dt} \left( \int_S \vec{B} \cdot d\vec{A} \right) \quad \left| \quad \oint_C \vec{H} \cdot d\vec{l} = \int_S \vec{J} \cdot d\vec{A} + \frac{d}{dt} \int_S \epsilon E dA \right.$$

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## Uniform Electromagnetic Waves

$E_y$  varies along the  $z$ -direction and  $E$  is constant in the two other directions

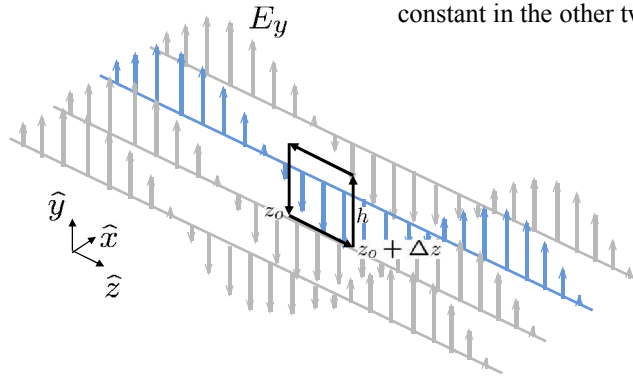


$$\oint_C \vec{E} \cdot d\vec{l} = -\frac{d}{dt} \left( \int_S \vec{B} \cdot d\vec{A} \right)$$

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## Uniform Electromagnetic Waves

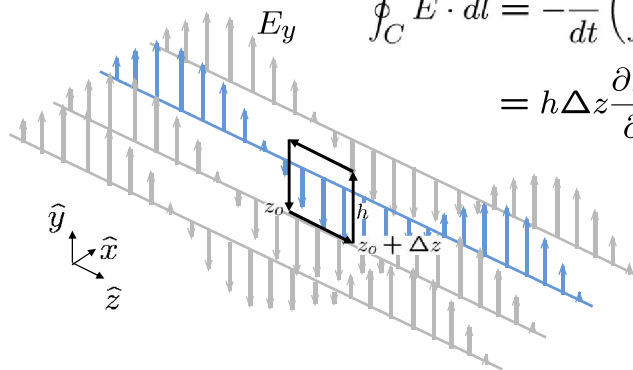
$E_y$  varies along the  $z$ -direction and  $E_y$  is constant in the other two directions



$$\begin{aligned} \oint_C \vec{E} \cdot d\vec{l} &= hE_y(z_0 + \Delta z) - hE_y(z_0) = -\frac{d}{dt} \left( \int_S \vec{B} \cdot d\vec{A} \right) \\ &= h\Delta z \frac{E_y(z_0 + \Delta z) - E_y(z_0)}{\Delta z} = h\Delta z \frac{\partial E_y}{\partial z} \end{aligned}$$

## Electromagnetic Waves

$$\begin{aligned} \oint_C \vec{E} \cdot d\vec{l} &= -\frac{d}{dt} \left( \int_S \vec{B} \cdot d\vec{A} \right) \\ &= h\Delta z \frac{\partial E_y}{\partial z} \end{aligned}$$



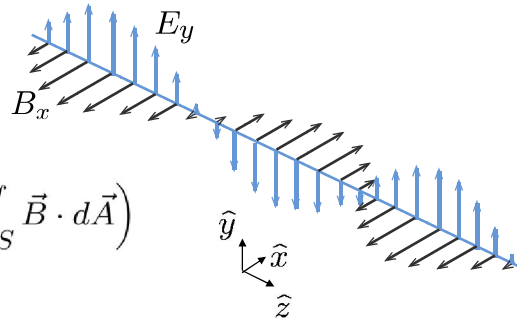
$E_y$ -field cannot vary in  $z$ -direction without a time-varying  $B$ -field ...

$$-\frac{\partial}{\partial t} \int \vec{B} d\vec{A} = -\frac{\partial}{\partial t} (-B_x h \Delta z) = h\Delta z \frac{\partial B_x}{\partial t}$$

...and waves must have both electric and magnetic components !

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## Uniform Electromagnetic Plane Waves



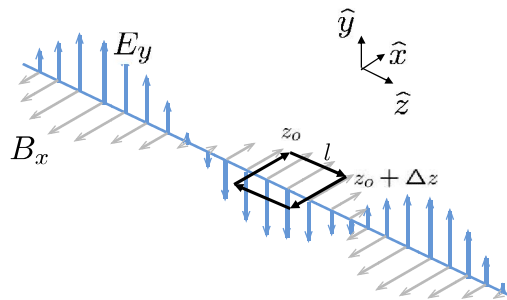
$$\oint_C \vec{E} \cdot d\vec{l} = -\frac{d}{dt} \left( \int_S \vec{B} \cdot d\vec{A} \right)$$

$$\frac{\partial E_y}{\partial z} = \frac{\partial B_x}{\partial t}$$

The **y-component of E** that varies across space is associated with the **x-component of B** that varies in time

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## Uniform Electromagnetic Plane Waves



$$\oint_C \vec{H} \cdot d\vec{l} = \int_S \vec{J} \cdot d\vec{A} + \frac{d}{dt} \int_S \epsilon_0 \vec{E} \cdot d\vec{A}$$

$$\frac{\partial B_x(z_0)}{\partial z} = \epsilon_0 \mu_0 \frac{\partial E_y}{\partial t} \quad \text{Source free: } J = 0$$

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## *The Wave Equation*

Time-varying  $E_y$  generates  
spatially varying  $B_z \dots$

$$\frac{\partial^2 B_x(z_0)}{\partial z \partial t} = \epsilon_0 \mu_0 \frac{\partial^2 E_y}{\partial t^2}$$

Time-varying  $B_z$  generates  
spatially varying  $E_y \dots$

$$\frac{\partial^2 E_y}{\partial z^2} = \frac{\partial^2 B_x(z_0)}{\partial t \partial z}$$

The temporal and spatial variations in  $E_y$  are coupled together to yield ...

$$\boxed{\frac{\partial^2 E_y}{\partial z^2} = \epsilon_0 \mu_0 \frac{\partial^2 E_y}{\partial t^2}}$$

... the Wave Equation.

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## *Wave Equation via Differential Equations*

$$\text{Faraday: } \underbrace{\begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ E_x & E_y & E_z \end{vmatrix}}_{\nabla \times \vec{E}} = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ 0 & E_y & 0 \end{vmatrix} = -\hat{x} \frac{\partial E_y}{\partial z} = -\hat{x} \frac{\partial \mu H_x}{\partial t}$$

$$\text{Ampere: } \underbrace{\begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ H_x & H_y & H_z \end{vmatrix}}_{\nabla \times \vec{H}} = \begin{vmatrix} \hat{x} & \hat{y} & \hat{z} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ H_x & 0 & 0 \end{vmatrix} = -\hat{y} \frac{\partial H_x}{\partial z} = -\hat{y} \frac{\partial \epsilon E_y}{\partial t}$$

Substitution yields the wave equation:

$$\frac{\partial^2 E_y}{\partial z^2} = \epsilon_0 \mu_0 \frac{\partial^2 E_y}{\partial t^2}$$

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## Uniform Plane Wave Solutions

The 1-D wave equation

$$\frac{\partial^2 E_y}{\partial z^2} - \epsilon_0 \mu_0 \frac{\partial^2 E_y}{\partial t^2} = 0$$

- $E_y(z,t)$  is any function for which the second derivative in space equals its second derivative in time, times a constant. The solution is therefore any function with the same dependence on time as on space, e.g.

$$E_y = f_+(t - z/c) + f_-(t + z/c)$$

- The functions  $f_+(z-ct)$  and  $f_-(z+ct)$  represent uniform waves propagating in the  $+z$  and  $-z$  directions respectively.

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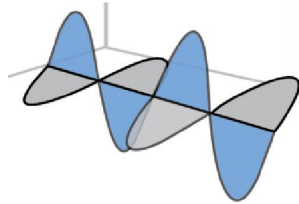
## Speed of Light

- The *velocity of propagation* is determined solely by the dielectric permittivity and magnetic permeability:

$$c = \frac{1}{\sqrt{\epsilon_0 \mu_0}}$$

- The functions  $f_+$  and  $f_-$  are determined by the source and the other boundary conditions.

$$E_y = f_+(t - z/c) + f_-(t + z/c)$$



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